

直边上具有混合支撑段薄矩形板的弯曲

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摘 要

在本文中, 应用功的互等定理法首次给出了直边上具有简支段与自由段混合支撑的矩形板的精确解析解。

作为比较, 我们用有限元法计算了同一问题。比较表明, 该解析解是正确的。

关键词 功的互等定理法 支撑段 三角级数的转换

一、引 言

在文章[1]中, 我们首先应用功的互等定理求解了悬臂矩形板的弯曲。之后, 我们继续推广功的互等定理的应用于求解板壳力学的平衡、振动和稳定, 于求解弹性力学问题的位移方程, 因此, 形成了一个系统的方法, 我们称其为功的互等定理法^[2~4]。

在本文中, 我们将应用功的互等定理法于求解在直边上具有混合支撑段的薄矩形板的挠曲方程并首先给出了在直边上具有简支段与自由段混合支撑的矩形板的精确解析解。

分析和计算表明, 功的互等定理法比叠加法要简便有效得多^[6]。

二、功的互等定理法

基本系统, 即在一单位集中荷载作用下的四边简支矩形板示于图1。图2所示为在均布荷载 q_0 作用下的实际系统。

我们假设, 实际系统直边自由段的挠度表达式分别为

$$W_{s,i}(y) = \sum_{s=1,2}^{\infty} E_{s,i} \sin \frac{s\pi(y-b_i)}{l_i} \quad (i=1,2,\dots,\alpha) \quad (2.1)$$

$$W_{s,i}(y) = \sum_{s=1,2}^{\infty} E_{s,i} \sin \frac{s\pi(y-b_i)}{l_i} \quad (i=1,2,\dots,\beta) \quad (2.2)$$

$$W_{s,i}(x) = \sum_{s=1,2}^{\infty} F_{s,i} \sin \frac{s\pi(x-a_i)}{h_i} \quad (i=1,2,\dots,\gamma) \quad (2.3)$$

$$W_{\bar{y}i}(x) = \sum_{s=1,2}^{\infty} \bar{F}_{s,i} \sin \frac{s\pi(x-\bar{a}_i)}{\bar{h}_i} \quad (i=1,2,\dots,\delta) \quad (2.4)$$

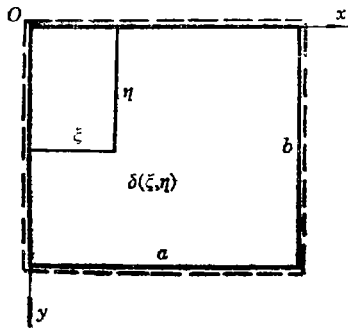


图 1

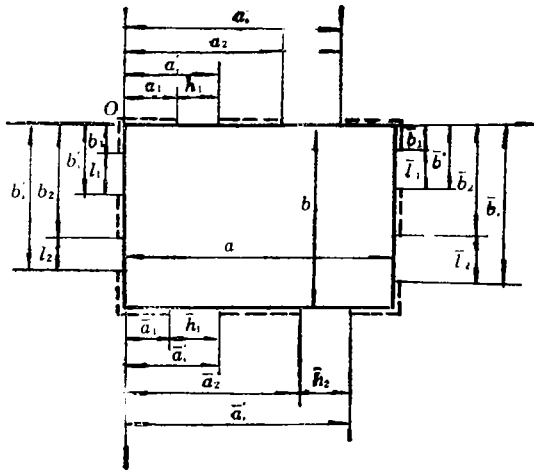


图 2

其中所有的符号示于图2。在基本系统与实际系统之间应用功的互等定理，我们得到

$$\begin{aligned} W(\xi, \eta) = & \int_0^a \int_0^b q_0 W_0(x, y, \xi, \eta) dx dy + \sum_{i=1,2}^{\alpha} \int_{b_i}^{b_i'} V_{0x}(x, y, \xi, \eta) \Big|_{x=0} W_{xi}(y) dy \\ & - \sum_{i=1,2}^{\beta} \int_{\bar{b}_i}^{\bar{b}_i'} V_{0x}(x, y, \xi, \eta) \Big|_{x=a} W_{xi}(y) dy \\ & + \sum_{i=1,2}^{\gamma} \int_{a_i}^{a_i'} V_{0y}(x, y, \xi, \eta) \Big|_{y=0} W_{yi}(x) dx \\ & - \sum_{i=1,2}^{\delta} \int_{\bar{a}_i}^{\bar{a}_i'} V_{0y}(x, y, \xi, \eta) \Big|_{y=b} W_{yi}(x) dx \end{aligned} \quad (2.5)$$

其中 $W_0(x, y, \rho, \eta)$ 为基本系统的基本解， V_{0x} 和 V_{0y} 是基本解的等效切力。

实际系统的边界条件为

$$V_{\xi}|_{\xi=0} = 0 \quad (b_i \leq \eta \leq b_i', i=1, 2, \dots, \alpha) \quad (2.6)$$

$$V_{\xi}|_{\xi=a} = 0 \quad (\bar{b}_i \leq \eta \leq \bar{b}_i', i=1, 2, \dots, \beta) \quad (2.7)$$

$$V_{\eta}|_{\eta=0} = 0 \quad (a_i \leq \xi \leq a_i', i=1, 2, \dots, \gamma) \quad (2.8)$$

$$V_{\eta}|_{\eta=b} = 0 \quad (\bar{a}_i \leq \xi \leq \bar{a}_i', i=1, 2, \dots, \delta) \quad (2.9)$$

剩下的简支段边界条件自动满足。

转换展开在矩形板每边的三角级数为展开在每一自由段的三角级数，我们得到

$$\begin{aligned} W(\xi, \eta) = & \frac{4q_0 b^4}{\pi^5 D} \sum_{n=1,3}^{\infty} \frac{1}{n^5} \left\{ 1 + \frac{1}{2 \operatorname{ch} \frac{\alpha_n}{2}} \left[\frac{\alpha_n}{a} \left(\xi - \frac{a}{2} \right) \operatorname{sh} \frac{\alpha_n}{a} \left(\xi - \frac{a}{2} \right) \right. \right. \\ & \left. \left. - \left(2 + \frac{\alpha_n}{2} \operatorname{th} \frac{\alpha_n}{2} \right) \operatorname{ch} \frac{\alpha_n}{a} \left(\xi - \frac{a}{2} \right) \right] \right\} \sin \frac{n\pi\eta}{b} \end{aligned}$$

$$\begin{aligned}
& - \frac{1}{\pi b} \sum_{i=1,2}^a l_i \sum_{s=1,2}^{\infty} s E_{s,i} \sum_{n=1,2}^{\infty} \left\{ 2 + (1-\nu) \left[\alpha_n \operatorname{cth} \alpha_n \right. \right. \\
& \left. \left. - \frac{\alpha_n(a-\xi)}{a} \operatorname{cth} \frac{\alpha_n(a-\xi)}{a} \right] \right\} \frac{1}{\operatorname{sh} \alpha_n} \operatorname{sh} \frac{\alpha_n(a-\xi)}{a} \sin \frac{n\pi\eta}{b} \\
& \cdot \Phi_{s,n}(b, l_i) - \frac{1}{\pi b} \sum_{i=1,2}^b \bar{l}_i \sum_{q=1,2}^{\infty} s \bar{E}_{s,i} \sum_{n=1,2}^{\infty} \left[2 + (1-\nu) \right. \\
& \left. \cdot \left(\alpha_n \operatorname{cth} \alpha_n - \frac{\alpha_n \xi}{a} \operatorname{cth} \frac{\alpha_n \xi}{a} \right) \right] \frac{1}{\operatorname{sh} \alpha_n} \operatorname{sh} \frac{\alpha_n \xi}{a} \sin \frac{n\pi\eta}{b} \Phi_{s,n}(b, l_i) \\
& - \frac{a^3}{b^4} \frac{4}{\pi^2} \sum_{i=1,2}^{\gamma} h_i \sum_{s=1,2}^{\infty} s F_{s,i} \sum_{m=1,2}^{\infty} \sum_{n=1,2}^{\infty} \frac{n \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \\
& \cdot \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b} \Phi_{s,m}(a, h_i) + \frac{a^3}{b^4} \frac{4}{\pi^2} \sum_{i=1,2}^{\delta} \bar{h}_i \\
& \cdot \sum_{s=1,2}^{\infty} s \bar{F}_{s,i} \sum_{m=1,2}^{\infty} \sum_{n=1,2}^{\infty} \frac{(-1)^n n \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \\
& \cdot \sin \frac{m\pi\xi}{a} \sin \frac{n\pi\eta}{b} \Phi_{s,m}(a, \bar{h}_i) \tag{2.10}
\end{aligned}$$

其中 $\alpha_n = \frac{n\pi a}{b}$ 和 $\beta_m = \frac{m\pi b}{a}$.

将式(2.10)代入式(2.6)~(2.9)并转换三角级数, 我们分别得到

$$\begin{aligned}
& \sum_{i=1,2}^a l_i \sum_{s=1,2}^{\infty} s E_{s,i} \sum_{n=1,2}^{\infty} n^3 \frac{(1-\nu)\alpha_n + \frac{3+\nu}{2} \operatorname{sh} 2\alpha_n}{\operatorname{sh}^2 \alpha_n} \Phi_{s,n}(b, l_i) \Phi_{k,n}(b, l_j) \\
& - \sum_{i=1,2}^b \bar{l}_i \sum_{s=1,2}^{\infty} s \bar{E}_{s,i} \sum_{n=1,2}^{\infty} n^3 \frac{(1-\nu)\alpha_n \operatorname{ch} \alpha_n + (3+\nu) \operatorname{sh} \alpha_n}{\operatorname{sh}^2 \alpha_n} \\
& \cdot \Phi_{s,n}(b, l_i) \Phi_{k,n}(b, l_j) + \frac{4}{\pi(\nu-1)} \sum_{i=1,2}^{\gamma} h_i \sum_{s=1,2}^{\infty} s F_{s,i} \\
& \cdot \sum_{n=1,2}^{\infty} \sum_{m=1,2}^{\infty} \frac{mn \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right] \left[m^2 + (2-\nu) \left(\frac{a}{b} n \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \Phi_{s,m}(a, h_i) \\
& \cdot \Phi_{k,n}(b, l_j) - \frac{4}{\pi(\nu-1)} \sum_{i=1,2}^{\delta} \bar{h}_i \sum_{s=1,2}^{\infty} s \bar{F}_{s,i}
\end{aligned}$$

$$\begin{aligned} & \cdot \sum_{n=1,2} \sum_{m=1,2}^{\infty} \frac{(-1)^n mn \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right] \left[m^2 + (2-\nu) \left(\frac{a}{b} n \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \\ & \cdot \Phi_{sm}(a, \bar{h}_i) \Phi_{kn}(b, l_j) - \frac{q_0 b^5}{D} \frac{1}{\pi^4 (\nu-1)} \sum_{n=1,2}^{\infty} \frac{\alpha_n (\nu-1) + (3-\nu) \text{sh} \alpha_n}{n^2 \text{ch}^2 \frac{\alpha_n}{2}} \\ & \cdot \Phi_{kn}(b, l_j) = 0 \quad (k=1, 2, \dots, \infty; j=1, 2, \dots, a) \end{aligned} \quad (2.11)$$

$$\begin{aligned} & \sum_{i=1,2}^a l_i \sum_{s=1,2}^{\infty} s E_{s,i} \sum_{n=1,2}^{\infty} n^3 \frac{(1-\nu) \alpha_n \text{ch} \alpha_n + (3+\nu) \text{sh} \alpha_n}{\text{sh}^2 \alpha_n} \Phi_{sn}(b, l_i) \Phi_{kn}(b, l_j) \\ & - \sum_{i=1,2}^{\beta} \bar{l}_i \sum_{s=1,2}^{\infty} s \bar{E}_{s,i} \sum_{n=1,2}^{\infty} n^3 \frac{(1-\nu) \alpha_n + \frac{3+\nu}{2} \text{sh} 2\alpha_n}{\text{sh}^2 \alpha_n} \Phi_{sn}(b, l_i) \Phi_{kn}(b, l_j) \\ & + \frac{4}{\pi(\nu-1)} \sum_{i=1,2}^{\gamma} h_i \sum_{s=1,2}^{\infty} s F_{s,i} \\ & \cdot \sum_{n=1,2}^{\infty} \sum_{m=1,2}^{\infty} \frac{(-1)^{m+n} mn \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right] \left[m^2 + (2-\nu) \left(\frac{a}{b} n \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \\ & \cdot \Phi_{sm}(a, h_i) \Phi_{kn}(b, l_j) - \frac{4}{\pi(\nu-1)} \sum_{i=1,2}^{\delta} \bar{h}_i \sum_{s=1,2}^{\infty} s \bar{F}_{s,i} \\ & \cdot \sum_{n=1,2}^{\infty} \sum_{m=1,2}^{\infty} \frac{(-1)^{m+n} mn \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right] \left[m^2 + (2-\nu) \left(\frac{b}{a} n \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \Phi_{sm}(a, \bar{h}_i) \Phi_{kn}(b, l_i) \\ & + \frac{q_0 b^5}{D} \frac{1}{\pi^4 (\nu-1)} \sum_{n=1,3}^{\infty} \frac{\alpha_n (\nu-1) + (3-\nu) \text{sh} \alpha_n}{n^2 \text{ch}^2 \frac{\alpha_n}{2}} \Phi_{kn}(b, l_j) = 0 \end{aligned} \quad (k=1, 2, \dots, \infty; j=1, 2, \dots, \beta) \quad (2.12)$$

$$\begin{aligned} & \sum_{i=1,2}^a l_i \sum_{s=1,2}^{\infty} s E_{s,i} \\ & \cdot \sum_{m=1,2}^{\infty} \sum_{n=1,2}^{\infty} \frac{mn \left[m^2 + (2-\nu) \left(\frac{a}{b} n \right)^2 \right] \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \Phi_{sn}(b, l_i) \Phi_{km}(a, h_j) \\ & - \sum_{i=1,2}^{\beta} \bar{l}_i \sum_{s=1,2}^{\infty} s \bar{E}_{s,i} \sum_{n=1,2}^{\infty} \sum_{m=1,2}^{\infty} \frac{(-1)^m mn \left[m^2 + (2-\nu) \left(\frac{a}{b} n \right)^2 \right] \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \\ & \cdot \Phi_{sn}(b, l_i) \Phi_{km}(a, l_j) + \left(\frac{b}{a} \right)^4 \frac{\pi(\nu-1)}{4} \sum_{i=1,2}^{\gamma} h_i \sum_{s=1,2}^{\infty} s F_{s,i} \end{aligned}$$

$$\begin{aligned}
& \cdot \sum_{m=1,2}^{\infty} m^3 \frac{(1-\nu)\beta_m + \frac{3+\nu}{2} \text{sh}^2 \beta_m}{\text{sh}^2 \beta_m} \Phi_{sm}(a, h_i) \Phi_{km}(a, h_j) \\
& - \left(\frac{b}{a}\right)^4 \frac{\pi(\nu-1)}{4} \sum_{i=1,2}^{\delta} \bar{h}_i \sum_{s=1,2}^{\infty} s \bar{F}_{s,i} \sum_{m=1,2}^{\infty} m^3 \frac{(1-\nu)\beta_m \text{ch} \beta_m + (3+\nu) \text{sh} \beta_m}{\text{sh}^2 \beta_m} \\
& \cdot \Phi_{sm}(a, \bar{h}_i) \Phi_{km}(a, h_j) - \frac{q_0 a b^4}{D} \frac{1}{4\pi^3} \sum_{m=1,3}^{\infty} \frac{\beta_m(\nu-1) + (3-\nu) \text{sh} \beta_m}{m^2 \text{ch}^2 \frac{\beta_m}{2}} \\
& \cdot \Phi_{km}(a, h_j) = 0 \quad (k=1,2,\dots,\infty; j=1,2,\dots,\nu) \quad (2.13) \\
& \sum_{i=1,2}^{\alpha} l_i \sum_{s=1,2}^{\infty} s E_{s,i} \\
& \cdot \sum_{n=1,2}^{\infty} \sum_{m=1,2}^{\infty} \frac{(-1)^n m n \left[m^2 + (2-\nu) \left(\frac{a}{b} n\right)^2 \right] \left[n^2 + (2-\nu) \left(\frac{b}{a} m\right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n\right)^2 \right]^2} \\
& \cdot \Phi_{sn}(b, l_i) \Phi_{km}(a, \bar{h}_j) - \sum_{i=1,2}^{\beta} l_i \sum_{s=1,2}^{\infty} s E_{s,i} \\
& \cdot \sum_{n=1,2}^{\infty} \sum_{m=1,2}^{\infty} \frac{(-1)^{m+n} m n \left[m^2 + (2-\nu) \left(\frac{a}{b} n\right)^2 \right] \left[n^2 + (2-\nu) \left(\frac{b}{a} m\right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n\right)^2 \right]^2} \\
& \cdot \Phi_{sn}(b, l_i) \Phi_{km}(a, \bar{h}_j) + \left(\frac{b}{a}\right)^4 \frac{\pi(\nu-1)}{4} \sum_{i=1,2}^{\nu} h_i \sum_{s=1,2}^{\infty} s F_{s,i} \\
& \cdot \sum_{m=1,2}^{\infty} m^3 \frac{(1-\nu)\beta_m \text{ch} \beta_m + (3+\nu) \text{sh} \beta_m}{\text{sh}^2 \beta_m} \Phi_{sm}(a, h_i) \Phi_{km}(a, \bar{h}_j) \\
& - \left(\frac{b}{a}\right)^4 \frac{\pi(\nu-1)}{4} \sum_{i=1,2}^{\delta} \bar{h}_i \sum_{s=1,2}^{\infty} s \bar{F}_{s,i} \\
& \cdot \sum_{m=1,2}^{\infty} m^3 \frac{(1-\nu)\beta_m + \frac{3+\nu}{2} \text{sh}^2 \beta_m}{\text{sh}^2 \beta_m} \Phi_{sm}(a, \bar{h}_i) \Phi_{km}(a, \bar{h}_j) \\
& + \frac{q_0 a b^4}{D} \frac{1}{4\pi^3} \sum_{m=1,3}^{\infty} \frac{\beta_m(\nu-1) + (3-\nu) \text{sh} \beta_m}{m^2 \text{ch}^2 \frac{\beta_m}{2}} \Phi_{km}(a, \bar{h}_j) = 0 \\
& \quad (k=1,2,\dots,\infty; j=1,2,\dots,\delta) \quad (2.14)
\end{aligned}$$

我们应当指出，为加快收敛速度，一些三角级数必须转换成双曲函数。

三、四边中部为自由段、其余部分为简支段的矩形板

图3所示为一在均布载荷作用下的实际系统。其中 $a_1 = \bar{a}_1$, $h_1 = \bar{h}_1$, $b_1 = \bar{b}_1$ 和 $l_1 = \bar{l}_1$, 因

此 该板的弯曲是对称的, 并且 $E_{s,1} = \bar{E}_{s,1}$ 和 $F_{s,1} = \bar{F}_{s,1}$. 我们假设

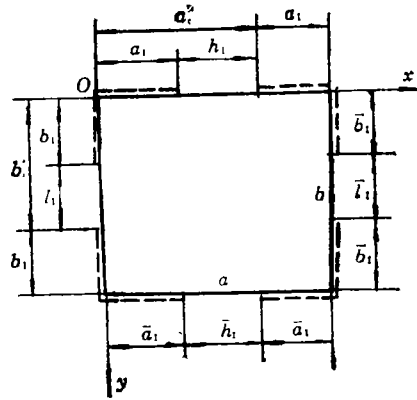


图 3

$$W_{s,1}(y) = W_{\bar{s},1}(y) = \sum_{s=1,3}^{\infty} E_{s,1} \sin \frac{s\pi(y-b)}{l_1} \quad (b_1 \leq y \leq b'_1) \quad (3.1)$$

$$W_{s,1}(x) = W_{\bar{s},1}(x) = \sum_{s=1,3}^{\infty} F_{s,1} \sin \frac{s\pi(x-a_1)}{h_1} \quad (a_1 \leq x \leq a'_1) \quad (3.2)$$

注意到

$$\Phi_{sm}(a, h_1) = -2 \frac{\sin \frac{m\pi}{a}}{s^2 - \left(\frac{h_1}{a} m\right)^2} \quad (\text{当 } m+s = \text{偶数}) \quad (3.3)$$

$$\Phi_{sm}(a, h_1) = 0 \quad (\text{当 } m+s = \text{奇数}) \quad (3.4)$$

我们能简化四个边界条件(2.11)~(2.14)为两个边界条件(2.12)和(2.13), 且它们分别为

$$\begin{aligned} & \sum_{s=1,3}^{\infty} s E_{s,1} \sum_{n=1,3}^{\infty} n^3 \frac{(1 - \text{ch} \alpha_n) [(1-\nu)\alpha_n - (3+\nu)\text{sh} \alpha_n]}{\text{sh}^2 \alpha_n} \\ & \cdot \frac{\sin^2 \frac{n\pi b_1}{b}}{\left[s^2 - \left(\frac{l_1}{b} n\right)^2\right] \left[k^2 - \left(\frac{l_1}{b} n\right)^2\right]} + \frac{8}{\pi(\nu-1)} \frac{h_1}{l_1} \\ & \cdot \sum_{s=1,3}^{\infty} s F_{s,1} \sum_{n=1,3}^{\infty} \sum_{m=1,3}^{\infty} \frac{mn \left[n^2 + (2-\nu) \left(\frac{b}{a} m\right)^2 \right] \left[m^2 + (2-\nu) \left(\frac{a}{b}\right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n\right)^2 \right]^2} \\ & \cdot \frac{\sin \frac{m\pi a_1}{a}}{s^2 - \left(\frac{h_1}{a} m\right)^2} \frac{\sin \frac{n\pi b_1}{b}}{k^2 - \left(\frac{l_1}{b} n\right)^2} = -\frac{q_0 b^5}{l_1 D} \frac{1}{2\pi^4(\nu-1)} \\ & \cdot \sum_{n=1,3}^{\infty} \frac{\alpha_n(\nu-1) + (3-\nu)\text{sh} \alpha_n}{n^2 \text{ch}^2 \frac{\alpha_n}{2}} \frac{\sin \frac{n\pi b_1}{b}}{k^2 - \left(\frac{l_1}{b} n\right)^2} \end{aligned} \quad (3.5)$$

$$\begin{aligned}
& \sum_{s=1,3}^{\infty} sE_{s,1} \cdot \sum_{m=1,3}^{\infty} \sum_{n=1,3}^{\infty} \frac{mn \left[m^2 + (2-\nu) \left(\frac{a}{b} n \right)^2 \right] \left[n^2 + (2-\nu) \left(\frac{b}{a} m \right)^2 \right]}{\left[m^2 + \left(\frac{a}{b} n \right)^2 \right]^2} \\
& \cdot \frac{\sin^2 \frac{n\pi b_1}{b}}{s^2 - \left(\frac{l_1}{b} n \right)^2} \frac{\sin^2 \frac{m\pi a_1}{a}}{k^2 - \left(\frac{h_1}{a} m \right)^2} + \frac{b^4}{a^4} \frac{\pi(\nu-1)}{8} \\
& \cdot \frac{h_1}{l_1} \sum_{s=1,3}^{\infty} sF_{s,1} \sum_{m=1,3}^{\infty} m^3 \frac{(1-\text{ch}\beta_m) [(1-\nu)\beta_m - (3+\nu)\text{sh}\beta_m]}{\text{sh}^2\beta_m} \\
& \cdot \frac{\sin^2 \frac{m\pi a_1}{a}}{\left[s^2 - \left(\frac{h_1}{a} m \right)^2 \right] \left[k^2 - \left(\frac{h_1}{a} m \right)^2 \right]} = - \frac{q_0 b^4 a}{l_1 D} \frac{1}{16\pi^3} \\
& \cdot \sum_{m=1,3}^{\infty} \frac{\beta_m(\nu-1) + (3-\nu)\text{sh}\beta_m}{m^2 \text{ch}^2 \frac{\beta_m}{2}} \frac{\sin^2 \frac{m\pi a_1}{a}}{k^2 - \left(\frac{h_1}{a} m \right)^2} \quad (3.6)
\end{aligned}$$

如果 $a=b$ 和 $l_1=h_1=1/3a$, 则有 $E_{s,1}=F_{s,1}$, (3.5) 和 (3.6) 是相互等价的, 且 (3.5) 成为

$$\begin{aligned}
& \sum_{s=1,3}^{\infty} sE_{s,1} \left\{ \sum_{n=1,3}^{\infty} n^3 \frac{(1-\text{ch}\alpha_n) [(1-\nu)\alpha_n - (3+\nu)\text{sh}\alpha_n]}{\text{sh}^2\alpha_n} \right. \\
& \cdot \left. \frac{\sin^2 \frac{n\pi b_1}{b}}{\left[s^2 - \left(\frac{l_1}{b} n \right)^2 \right] \left[k^2 - \left(\frac{l_1}{b} n \right)^2 \right]} + \frac{\pi(\nu-1)}{8} \right. \\
& \cdot \left. \sum_{n=1,3}^{\infty} \sum_{m=1,3}^{\infty} \frac{mn \left[n^2 + (2-\nu)m^2 \right] \left[m^2 + (2-\nu)n^2 \right]}{(m^2+n^2)^2} \cdot \frac{\sin^2 \frac{n\pi b_1}{b}}{\left[s^2 - \left(\frac{l_1}{b} m \right)^2 \right] \left[k^2 - \left(\frac{l_1}{b} n \right)^2 \right]} \right\} \\
& = - \frac{3q_0 b^4}{D} \frac{1}{2\pi^4(\nu-1)} \sum_{n=1,3}^{\infty} \frac{\alpha_n(\nu-1) + (3-\nu)\text{sh}\alpha_n}{n^2 \text{ch}^2 \frac{\alpha_n}{2}} \frac{\sin^2 \frac{n\pi b_1}{b}}{k^2 - \left(\frac{l_1}{b} n \right)^2} \\
& \quad (k=1, 3, \dots, \infty) \quad (3.7)
\end{aligned}$$

由附录, 三角级数到双曲函数的转换, 式(3.7)转换为

$$\begin{aligned}
& \sum_{s=1,3}^{\infty} sE_{s,1} \left\{ \sum_{n=1,3}^{\infty} n^3 \left[\frac{\alpha_n}{2} (1-\nu) \left(1 - \text{th}^2 \frac{\alpha_n}{2} \right) - (3+\nu) \text{th} \frac{\alpha_n}{2} \right] \right. \\
& \cdot \left. \frac{\sin^2 \frac{n\pi b_1}{b}}{\left[s^2 - \left(\frac{l_1}{b} n \right)^2 \right] \left[k^2 - \left(\frac{l_1}{b} n \right)^2 \right]} \right\}^{-2(1-\nu)}
\end{aligned}$$

$$\begin{aligned}
& \sum_{m=1,3}^{\infty} \frac{m^4}{k^2 + \left(\frac{l_1}{b} m\right)^2} \sin \frac{m\pi b_1}{b} \left\{ \frac{m \left(\frac{l_1}{b}\right)^2}{k^2 + \left(\frac{l_1}{b} m\right)^2} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} \right. \\
& \left. + \frac{\pi}{4} \left[\operatorname{th} \frac{m\pi}{2} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} - \frac{l_1}{b} \frac{\operatorname{sh} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} \right] - \frac{1}{m} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} \right\} \\
& = \frac{3q_0 b^4}{D} \frac{1}{2\pi^4(\nu-1)} \sum_{n=1,3}^{\infty} \frac{1}{n^2} \left[\alpha_n(\nu-1) \left(1 - \operatorname{th}^2 \frac{\alpha_n}{2}\right) \right. \\
& \left. + 2(3-\nu) \operatorname{th} \frac{\alpha_n}{2} \right] \frac{\sin \frac{n\pi b_1}{b}}{k^2 - \left(\frac{l_1}{b} n\right)^2} \quad (k=1,3,\dots,\infty) \quad (3.8)
\end{aligned}$$

当 $s:n=l_1:b=1:3$, $k:n=l_1:b=1:3$ 和 $s:m=l_1:b=1:3$ 时, 方程(3.8)是可解的. 对 k 和 s 都取前十二项, 且在 VAX-780 计算机上计算, 我们得表 1~表 4.

图 1 自由段的挠度系数 $\left(\frac{q_0 a^4}{D} \times 10^{-3}\right)$

$E_{1,1}$	$E_{3,1}$	$E_{5,1}$	$E_{7,1}$	$E_{9,1}$	$E_{11,1}$
0.27949	-0.015169	-0.0053437	-0.0023878	-0.0012499	-0.00072258
$E_{13,1}$	$E_{15,1}$	$E_{17,1}$	$E_{19,1}$	$E_{21,1}$	$E_{23,1}$
-0.00044555	-0.00028643	-0.00018873	-0.00012553	-0.000082880	-0.000053021

表 2 $\nu=0$ 自由段的挠度分布 $\left(\frac{q_0 a^4}{D} \times 10^{-3}\right)$

方法 \ x/a	9/24	10/24	11/24	12/24
本文	0.0886	0.192	0.265	0.291
有限元(144单元)	0.083	0.185	0.257	0.283
误差	-6.3%	-3.6%	-3.0%	-2.7%

表 3 $x=\frac{a}{2}$ 的挠度分布 $\left(10^{-2} \times \frac{q_0 a^4}{D}\right)$

方法 \ y/b	2/12	2/12	3/12	4/12	5/12	6/12
本文	0.1297	0.2263	0.3081	0.3697	0.4078	0.4207
有限元(144单元)	0.1296	0.2256	0.3073	0.3687	0.4068	0.4196
误差	-0.77%	-0.31%	-0.26%	-0.27%	-0.25%	-0.26%

表 4

 $x = \frac{a}{2}$ 的弯矩分布 ($10^{-1} \times q_0 a^2$)

y/b	类别	M_x			M_y		
	方法	本 文	有限元(144单元)	误 差	本 文	有限元(144单元)	误 差
1/12		0.2438	0.2701	10.8%	0.1587	0.1529	-3.7%
2/12		0.3099	0.3156	1.8%	0.2905	0.2889	-0.55%
3/12		0.3807	0.3811	0.11%	0.3818	0.3818	0
4/12		0.4374	0.4369	-0.11%	0.4404	0.4403	-0.023%
5/12		0.4726	0.4721	-0.11%	0.4736	0.4731	-0.11%
6/12		0.4844	0.4839	-0.10%	0.4844	0.4837	-0.14%

板单元的划分示于图4。

四、结 论

由分析和计算我们可得出下述结论：对于求解直边上有混合支撑段矩形板的挠曲方程，

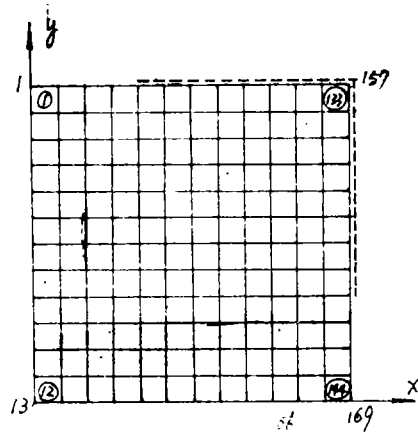


图 4

功的互等定理法是简单、方便和有效的。

必须指出，本法比国内外长期广泛采用的叠加法要简单的多。

附 录

在方程(3.7)左边第二项中，我们假设

$$A = (2-\nu)m^6 A_1 + [1+(2-\nu)^2]m^3 A_2 + (2-\nu)m A_3 \quad (\text{A.1})$$

$$\text{其中 } A_1 = \sum_{n=1,3}^{\infty} \frac{n \sin \frac{n\pi b_1}{b}}{(m^2+n^2)^2 \left[k^2 - \left(\frac{l_1}{b} n \right)^2 \right]} \quad (\text{A.2})$$

$$A_2 = \sum_{n=1,3}^{\infty} \frac{n^3 \sin \frac{n\pi b_1}{b}}{(m^2 + n^2)^2 \left[k^2 - \left(\frac{l_1}{b} n \right)^2 \right]} \quad (\text{A.3})$$

$$A_3 = \sum_{n=1,3}^{\infty} \frac{n^5 \sin \frac{n\pi b_1}{b}}{(m^2 + n^2)^2 \left[k^2 - \left(\frac{l_1}{b} n \right)^2 \right]} \quad (\text{A.4})$$

利用和[5]中(D1)~(D7)相同的推导, 我们得

$$A_1 = \frac{\pi}{4m} \frac{1}{k^2 + \left(\frac{l_1}{b} m \right)^2} \left\{ \frac{m \left(\frac{l_1}{b} \right)^2}{k^2 + \left(\frac{l_1}{b} m \right)^2} \frac{\text{ch} \frac{m\pi l_1}{2b}}{\text{ch} \frac{m\pi}{2}} \right. \\ \left. + \frac{\pi}{4} \left[\text{th} \frac{m\pi}{2} \frac{\text{ch} \frac{m\pi l_1}{2b}}{\text{ch} \frac{m\pi}{2}} - \frac{l_1}{b} \frac{\text{sh} \frac{m\pi l_1}{2b}}{\text{ch} \frac{m\pi}{2}} \right] \right\} \quad (\text{A.5})$$

利用和[5]中(D1)相似的方法, 我们得到

$$A_2 = - \frac{\left(\frac{b}{l_1} \right)^2}{\left[m^2 + \left(\frac{b}{l_1} k \right)^2 \right]^2} \left[\sum_{n=1,3}^{\infty} \frac{n^3}{n^4 - \left(\frac{b}{l_1} k \right)^2} \sin \frac{n\pi b_1}{b} \right. \\ \left. - \sum_{n=1,3}^{\infty} \frac{n^3}{n^2 + m^2} \sin \frac{n\pi b_1}{b} \right] + \frac{\left(\frac{b}{l_1} \right)^2}{m^2 + \left(\frac{b}{l_1} k \right)^2} \sum_{n=2,3}^{\infty} \frac{n^3}{(n^2 + m^2)^2} \sin \frac{n\pi b_1}{b} \quad (\text{A.6})$$

注意到

$$\sum_{n=1,3}^{\infty} \frac{n^3}{n^2 - \left(\frac{b}{l_1} k \right)^2} \sin \frac{n\pi b_1}{b} = \frac{\pi}{4} \left(\frac{b}{l_1} k \right)^2 \frac{\cos \frac{\pi}{2} k}{\cos \frac{b}{2l_1} k} = 0 \quad (k \text{ 是奇数}) \quad (\text{A.7})$$

$$\sum_{n=1,3}^{\infty} \frac{n^3}{n^2 + m^2} \sin \frac{n\pi b_1}{b} = - \frac{\pi}{4} m^2 \frac{\text{ch} \frac{m\pi b_1}{2b}}{\text{ch} \frac{m\pi}{2}} \quad (\text{A.8})$$

和

$$\sum_{n=1,3}^{\infty} \frac{n^3}{(n^2 + m^2)^2} \sin \frac{n\pi b_1}{b} = - \frac{m\pi^2}{16 \text{ch}^2 \frac{m\pi}{2}} \left(\text{sh} \frac{m\pi}{2} \text{ch} \frac{m\pi l_1}{2b} \right. \\ \left. - \frac{l_1}{b} \text{sh} \frac{m\pi l_1}{2b} \text{ch} \frac{m\pi}{2} - \frac{4}{m\pi} \text{ch} \frac{m\pi}{2} \text{ch} \frac{m\pi l_1}{2b} \right) \quad (\text{A.9})$$

并且将它们代入到(A.6)中, 我们得到

$$A_2 = - \frac{\pi}{4} m \frac{1}{k^2 + \left(\frac{l_1}{b} m \right)^2} \left\{ \frac{m \left(\frac{l_1}{b} \right)^2}{k^2 + \left(\frac{l_1}{b} m \right)^2} \frac{\text{ch} \frac{m\pi l_1}{2b}}{\text{ch} \frac{m\pi}{2}} \right. \\ \left. + \frac{\pi}{4} \left[\text{th} \frac{m\pi}{2} \frac{\text{ch} \frac{m\pi l_1}{2b}}{\text{ch} \frac{m\pi}{2}} - \frac{l_1}{b} \frac{\text{sh} \frac{m\pi l_1}{2b}}{\text{ch} \frac{m\pi}{2}} - \frac{4}{m\pi} \frac{\text{ch} \frac{m\pi l_1}{2b}}{\text{ch} \frac{m\pi}{2}} \right] \right\} \quad (\text{A.10})$$

用和导出 A_2 的相同方法, 我们得

$$A_3 = \frac{\pi}{4} m^3 \frac{1}{k^2 + \left(\frac{l_1}{b} m\right)^2} \left\{ \frac{m \left(\frac{l_1}{b}\right)^2}{k^2 + \left(\frac{l_1}{b} m\right)^2} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} + \frac{\pi}{4} \left[\operatorname{th} \frac{m\pi}{2} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} - \frac{l_1}{b} \frac{\operatorname{sh} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} - \frac{4}{m\pi} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} \right] \right\} \quad (\text{A.11})$$

将(A.5), (A.10)和(A.11)代入(A.1)中, 在最后得到

$$A = -\frac{\pi}{4} m^4 \frac{(1-\nu)^2}{k^2 + \left(\frac{l_1}{b} m\right)^2} \left\{ \frac{m \left(\frac{l_1}{b}\right)^2}{k^2 + \left(\frac{l_1}{b} m\right)^2} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} + \frac{\pi}{4} \left(\operatorname{th} \frac{m\pi}{2} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} - \frac{l_1}{b} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} - \frac{4}{m\pi} \frac{\operatorname{ch} \frac{m\pi l_1}{2b}}{\operatorname{ch} \frac{m\pi}{2}} \right) \right\} \quad (\text{A.12})$$

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The Bending of Thin Rectangular Plates with Mixed Supported Segments of Straight Edges

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Abstract

In this paper, the exact analytical solution of the rectangular plate having simply supported segments mixed with free segments of straight edges are first given by means of the method of reciprocal theorem.

By comparison, we calculate the same problem by finite element method. The comparison shows that the analytical solution is correct.

Key words the method of reciprocal theorem, supported segment, transformation of trigonometric series