

文章编号: 1000-0887(1999) 09\_0985\_06

# 关于层合板层间的连续条件\*

张慎学

(吉林大学数学系, 长春 130023)

(钱伟长推荐)

摘要: 给出了保证满足复合材料层合板层间位移和应力连续条件的各种形式混合能量原理的泛函, 并通过算例证明了泛函的可靠性。

关键词: 层合板; 层间连续条件; 混合能量; 泛函

中图分类号: O343.1 文献标识码: A

## 前 言

目前, 在层合板的理论研究和实际计算中, 人们所关注的一个重要问题之一, 是层间界面的位移和应力的连续条件。本文从这一关注出发, 给出了保证满足复合材料层合板层间位移和应力连续条件的各种形式混合能量原理的泛函, 并运用其中之一, 通过有限元数值计算证实了泛函的可靠性。

## 1 混合应力应变关系

设层合板的各层分别由均匀材料所组成, 应力应变关系记为

$$\sigma_i = Q_{ij} \varepsilon_j \quad (i, j = 1, 2, \dots, 6), \quad (1)$$

或

$$\varepsilon_i = q_{ij} \sigma_j \quad (i, j = 1, 2, \dots, 6), \quad (2)$$

式中

$$\left. \begin{aligned} \sigma_1 &= \sigma_{11}, & \sigma_2 &= \sigma_{22}, & \sigma_3 &= \sigma_{33}, & \sigma_4 &= \sigma_{23}, & \sigma_5 &= \sigma_{13}, & \sigma_6 &= \sigma_{12}, \\ \varepsilon_1 &= \varepsilon_{11}, & \varepsilon_2 &= \varepsilon_{22}, & \varepsilon_3 &= \varepsilon_{33}, & \varepsilon_4 &= 2\varepsilon_{23}, & \varepsilon_5 &= 2\varepsilon_{13}, & \varepsilon_6 &= 2\varepsilon_{12}, \end{aligned} \right\} \quad (3a, b)$$

$\sigma_{\lambda\mu}, \varepsilon_{\lambda\mu} (\lambda, \mu = 1, 2, 3)$  是应力和应变张量在直角笛卡儿坐标系  $(X_\lambda)$  下的分量。

将(3)的左边的次序重排, 记为

$$\Sigma = [ \sigma_p, \dots, \sigma_q; \sigma_r, \dots, \sigma_t ]^T = \left\{ \begin{array}{c} \Sigma_K \\ E_J \end{array} \right\} \quad (4)$$
$$E = [ \varepsilon_p, \dots, \varepsilon_q; \varepsilon_r, \dots, \varepsilon_t ]^T = \left\{ \begin{array}{c} \Sigma_K \\ E_J \end{array} \right\}$$

\* 收稿日期: 1997\_12\_21; 修订日期: 1999\_03\_12

作者简介: 张慎学(1940~), 男, 教授, 从事微分方程极小化与非奇异线性变换分解的研究, 发表论文 27 篇, 有 19 篇次被国内外检索刊物收入。

其中  $\Sigma_K = [\sigma_p, \dots, \sigma_q]^T$ ,  $\Sigma_J = [\sigma_r, \dots, \sigma_t]^T$  分别含  $K, J$  个分量,  $E_K = [\varepsilon_p, \dots, \varepsilon_q]^T$ ,  $E_J = [\varepsilon_r, \dots, \varepsilon_t]^T$  分别含  $K, J$  个分量, 且  $K + J = 6$ ,

如果我们写

$$\Sigma = RE, \quad E = r\Sigma, \tag{5}$$

则

$$R = \begin{bmatrix} Q_{pp} & \dots & Q_{pq} & Q_{pr} & \dots & Q_{pt} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ Q_{qp} & \dots & Q_{qq} & Q_{qr} & \dots & Q_{qt} \\ Q_{rp} & \dots & Q_{rq} & Q_{rr} & \dots & Q_{rt} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ Q_{tp} & \dots & Q_{tq} & Q_{tr} & \dots & Q_{tt} \end{bmatrix}, \quad r = \begin{bmatrix} q_{pp} & \dots & q_{pq} & q_{pr} & \dots & q_{pt} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ q_{qp} & \dots & q_{qq} & q_{qr} & \dots & q_{qt} \\ q_{rp} & \dots & q_{rq} & q_{rr} & \dots & q_{rt} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ q_{tp} & \dots & q_{tq} & q_{tr} & \dots & q_{tt} \end{bmatrix}, \tag{6}$$

是  $6 \times 6$  矩阵,  $Rr = I$  ( $I$  是单位矩阵)。

如果将 (5) 写成

$$\left. \begin{aligned} \left. \begin{aligned} \Sigma_K \\ \Sigma_J \end{aligned} \right\} &= \begin{bmatrix} R_{11} & R_{12} \\ R_{21} & R_{22} \end{bmatrix} \begin{bmatrix} E_K \\ E_J \end{bmatrix}, \\ \left. \begin{aligned} E_K \\ E_J \end{aligned} \right\} &= \begin{bmatrix} r_{11} & r_{12} \\ r_{21} & r_{22} \end{bmatrix} \begin{bmatrix} \Sigma_K \\ \Sigma_J \end{bmatrix}, \end{aligned} \right\} \tag{7}$$

则

$$\begin{pmatrix} \Sigma_K \\ \Sigma_J \end{pmatrix} = \begin{bmatrix} R_{11}^{-1} & R_{11}^{-1}R_{12} \\ R_{21}R_{11}^{-1} & R_{21}R_{11}^{-1}R_{12} - R_{22} \end{bmatrix} \begin{pmatrix} \Sigma_K \\ -E_J \end{pmatrix} = \begin{bmatrix} r_{11} - r_{12}r_{22}^{-1}r_{21} & -r_{12}r_{22}^{-1} \\ -r_{22}^{-1}r_{21} & -r_{22}^{-1} \end{bmatrix} \begin{pmatrix} \Sigma_K \\ -E_J \end{pmatrix}, \tag{8}$$

$$\begin{pmatrix} \Sigma_K \\ -E_J \end{pmatrix} = \begin{bmatrix} R_{11} - R_{12}R_{22}^{-1}R_{21} & R_{12}R_{22}^{-1} \\ R_{22}^{-1}R_{21} & -R_{22}^{-1} \end{bmatrix} \begin{pmatrix} E_K \\ \Sigma_J \end{pmatrix} = \begin{bmatrix} r_{11}^{-1} & -r_{11}^{-1}r_{12} \\ -r_{21}r_{11}^{-1} & r_{21}r_{11}^{-1}r_{12} - r_{22} \end{bmatrix} \begin{pmatrix} E_K \\ \Sigma_J \end{pmatrix}. \tag{9}$$

对于 (8), (9) 我们进一步用分量形式分别记成

$$s_i = C_{ij} r_j \tag{10}$$

$$r_i = H_{ij} s_j \tag{11}$$

其中  $C_{ij} = C_{ji}, H_{ij} = H_{ji}$  和

$$\left. \begin{aligned} (s_1, s_2, s_3, s_4, s_5, s_6) &= (\varepsilon_p, \dots, \varepsilon_q; \sigma_r, \dots, \sigma_t), \\ (r_1, r_2, r_3, r_4, r_5, r_6) &= (Q_p, \dots, Q_q; -\varepsilon_r, \dots, \varepsilon_t), \end{aligned} \right\} \tag{12}$$

我们定义混合能量密度

$$\overset{c}{W}_{KJ} = \frac{1}{2} C_{ij} r_i r_j, \quad \overset{c}{W}_{KJ} = \frac{1}{2} H_{ij} s_i s_j, \tag{13}$$

根据 (10), (11), (13), 我们得

$$\overset{c}{W}_{KJ, r_i} = s_i, \quad \overset{c}{W}_{KJ, s_i} = r_i \tag{14}$$

其中  $\overset{c}{W}_{KJ, r_i}$  和  $\overset{c}{W}_{KJ, s_i}$  表示  $\overset{c}{W}_{KJ}$  和  $\overset{c}{W}_{KJ}$  分别关于  $r_i$  和  $s_i$  的偏导数, 由此, 我们找到两类应力应变关

系•

$$\left. \begin{aligned} \bar{W}_{KJ, \sigma_\alpha} &= \varepsilon_\alpha \quad (\alpha = p, \dots, q), \\ \bar{W}_{KJ, \varepsilon_\beta} &= -\sigma_\beta \quad (\beta = r, \dots, t), \end{aligned} \right\} \quad (15)$$

$$\left. \begin{aligned} \bar{W}_{KJ, \varepsilon_\alpha} &= \sigma_\alpha \quad (\alpha = p, \dots, q), \\ \bar{W}_{KJ, \sigma_\beta} &= -\varepsilon_\beta \quad (\beta = r, \dots, t). \end{aligned} \right\} \quad (16)$$

对任意变分  $\delta\sigma_\alpha (\alpha = p, \dots, q), \delta\varepsilon_\beta (\beta = r, \dots, t)$ 

可以得到

$$\delta \bar{W}_{KJ} = \sum_{\alpha=p, \dots, q} \varepsilon_\alpha \delta \sigma_\alpha - \sum_{\beta=r, \dots, t} \sigma_\beta \delta \varepsilon_\beta \quad (17)$$

对任意变分  $\delta\varepsilon_\alpha (\alpha = p, \dots, q), \delta\sigma_\beta (\beta = r, \dots, t)$ 

可以得到

$$\delta W_{KJ} = \sum_{\alpha=p, \dots, q} \sigma_\alpha \delta \varepsilon_\alpha - \sum_{\beta=r, \dots, t} \varepsilon_\beta \delta \sigma_\beta \quad (18)$$

当应力应变关系被满足时, 我们有

$$\bar{W}_{KJ} = W_{KJ} \quad (19)$$

如果  $\alpha = 1, 2, 3, 6; \beta = 4, 5$  则正交各向异性和横观各向同性材料矩阵  $[C_{ij}], [q_{ij}]$  分别是

$$\left[ \begin{array}{cccccc} S_{11} & S_{12} & S_{13} & 0 & 0 & 0 \\ S_{21} & S_{22} & S_{23} & 0 & 0 & 0 \\ S_{31} & S_{32} & S_{33} & 0 & 0 & 0 \\ 0 & 0 & 0 & S_{66} & 0 & 0 \\ 0 & 0 & 0 & 0 & -1/S_{44} & 0 \\ 0 & 0 & 0 & 0 & 0 & -1/S_{55} \end{array} \right]$$

和

$$\left[ \begin{array}{cccccc} 1/E_{11} & -\nu_{21}/E_{22} & -\nu_{21}/E_{22} & 0 & 0 & 0 \\ -\nu_{21}/E_{22} & 1/E_{11} & -\nu_{23}/E_{22} & 0 & 0 & 0 \\ -\nu_{21}/E_{22} & -\nu_{23}/E_{22} & 1/E_{22} & 0 & 0 & 0 \\ 0 & 0 & 0 & 1/\mu_{12} & 0 & 0 \\ 0 & 0 & 0 & 0 & -1/\mu_{23} & 0 \\ 0 & 0 & 0 & 0 & 0 & -1/\mu_{12} \end{array} \right]$$

这里  $S_{ij}$  是弹性柔度系数,  $E_{11}, \dots, \mu_{12}$  是工程弹性常数<sup>[1]</sup>。

## 2 关于层合板的混合能量原理

设

$$\begin{aligned} \Gamma(u, \lambda, \sigma_i, \varepsilon, \dots, \varepsilon_t) = & \sum_{m=1}^N \int_{V_m} [-\bar{W}_{KJ} - \sum_{\beta=r, \dots, t} \varepsilon_\beta \sigma_\beta + \varepsilon_i(u) \sigma_i - \\ & u \lambda f - \lambda] dV - \frac{1}{2} \sum_m \int_{S_m} [u \, d\lambda - u \, u_\lambda] [p \, d\lambda - p \, u_\lambda] ds - \end{aligned}$$

$$\int_{S_0} p_{\lambda\lambda} ds - \int_{S_u} (u_{\lambda\lambda} - u_{\lambda}) p_{\lambda} ds, \quad (20)$$

$$\begin{aligned} \Pi(u_{\lambda}, \alpha_i, \varepsilon_p, \dots, \varepsilon_q) = & \sum_{m=1}^N \int_{V_m} \left\{ \sum_{\alpha=p, \dots, q}^{(m)} \sigma_{\alpha} \varepsilon_{\alpha} - w_{KL} + u_{\lambda} [ (F_{\lambda\lambda}, \sigma_{\lambda\mu}), \mu ] + \right. \\ & \left. f_{\lambda} \right\} dV - \frac{1}{2} \sum_m \int_{S_m} [ p_{d\lambda} + p_{u\lambda} ] [ u_{d\lambda} + u_{u\lambda} ] ds - \\ & \int_{S_u} u_{\lambda} p_{\lambda} ds - \int_{S_0} (p_{\lambda\lambda} - p_{\lambda}) u_{\lambda} ds, \end{aligned} \quad (21)$$

其中  $u_{\lambda}$  和  $f_{\lambda}$  分别是位移和体力分量,

$$p_{\lambda} = F_{\lambda\lambda} \sigma_{\lambda\mu} n_{\mu}, \quad F_{\lambda\lambda} = \delta_{\lambda\lambda} + u_{\lambda, \lambda} \quad (22)$$

$n_{\mu}$  是层合板表面向外单位法向量;  $W_{KJ}$ ,  $\varepsilon_{\beta}$  等是第  $m$  层在直角笛卡儿坐标系  $\{X_{\lambda}\}$  ( $X_s$  轴向下) 中所占区域  $V_m$  上  $W_{KJ}$  和  $\varepsilon_{\beta}$  等的取值;  $S_{1-M}$  和  $S_N$  分别是层合板的上下表面,  $S_m$  是  $V_m$  和  $V_{m+1}$  的分界面,  $S_{um}$  是  $V_m$  的上表面,  $S_{dm}$  是  $V_m$  的下表面, 即  $S_m = S_{dm} = S_{u(m+1)}$ ;  $u_{d\lambda}$ ,  $u_{u\lambda}$  等的定义域是  $S_{dm} = S_{u(m+1)}$ , 层合板共有  $M+N$  层, 中面 ( $X_s = 0$ ) 以上有  $M$  层;  $u_{\lambda}, p_{\lambda}$ , 等表示已知量

$$\varepsilon_{\lambda\lambda}(u) = \frac{1}{2} (u_{\lambda, \lambda} + u_{\lambda, \lambda} + u_{\mu, \mu\lambda}, \lambda), \quad (23)$$

其中  $( )$ ,  $\lambda$  表示  $( )$  关于  $X_{\lambda}$  的偏导数.

有时, 为了简化记号, 我们略去上标  $(m)$ , 例如: 式(1), (8) 等即是.

为了方便, 我们引进如下记号:

$$\begin{aligned} T_J &= (u_{\lambda}, \sigma_i, \varepsilon_r, \dots, \varepsilon_t) \in L, \text{ 当且仅当 } u_{\lambda} \in C^1(V_m), \\ \sigma_i &\in C^1(V_m), \varepsilon_r, \dots, \varepsilon_q \in C(V_m), U_K = (u_{\lambda}, \sigma_i, \\ \varepsilon_p, \dots, \varepsilon_q) &\in H_K \text{ 当且仅当 } u_{\lambda} \in C^1(V_m), \sigma_i \in C^1(V_m), \\ \varepsilon_r, \dots, \varepsilon_q &\in C(V_m). \end{aligned}$$

当  $T_J, T_J + \delta T_J \in L_J$  时, 我们得

$$\begin{aligned} \delta \Gamma = & \sum_{m=1}^N \int_{V_m} \left\{ \sum_{\alpha=p, \dots, q} [ \varepsilon_{\alpha}(u) - W_{KJ}, \sigma^{\alpha} ] \delta \sigma_{\alpha} - \right. \\ & \sum_{\beta=r, \dots, t} [ (W_{KJ}, \varepsilon_{\beta} + \alpha_{\beta}) \delta \varepsilon_{\beta} + (\varepsilon_{\beta} - \varepsilon_{\beta}(u)) \delta \alpha_{\beta} ] - \\ & [ (F_{\lambda\lambda}, \sigma_{\lambda\mu}), \mu + f_{\lambda} ] \delta u_{\lambda} \left. \right\} dV + \frac{1}{2} \sum_m \int_{S_m} \left\{ [ p_{d\lambda} + p_{u\lambda} ] \times \right. \\ & \left. [ u_{d\lambda} + u_{u\lambda} ] - [ u_{d\lambda} - u_{u\lambda} ] \delta [ p_{d\lambda} - p_{u\lambda} ] \right\} dS + \\ & \int_{S_0} (p_{\lambda\lambda} - p_{\lambda}) u_{\lambda} ds - \int_{S_u} (u_{\lambda\lambda} - u_{\lambda}) \delta p_{\lambda} ds \end{aligned} \quad (24)$$

当  $U_K, U_K + \delta K_K \in H_K$  时, 我们得

$$\delta \Pi = \sum_{m=1}^N \int_{V_m} \left\{ \sum_{\alpha=p, \dots, q} (\alpha_{\alpha} - W_{KJ}, \varepsilon_{\alpha}) \delta \varepsilon_{\alpha} - \right.$$

$$\begin{aligned} & \sum_{\beta=r, \dots, t} [(W_{KJ}, \sigma_{\beta+} \mathfrak{E}_{\beta}(u) \delta\sigma_{\beta+} + \sum_{\alpha=p, \dots, q} [\varepsilon_{\alpha} - \\ & \varepsilon_{\alpha}(u)] \delta\sigma_{\alpha+} + [(F_{\lambda\lambda}, \sigma_{\lambda\mu}), \mu + f_{\lambda}] \delta u_{\lambda}] dV - \frac{1}{2} \sum_m \int_{S_m} \left\{ [p \ d\lambda + \right. \\ & \left. p \ u_{\lambda}] \delta [u \ d\lambda + u \ u_{\lambda}] + [u \ u_{\lambda-} \ u \ d\lambda] \delta [p \ d\lambda - p \ u_{\lambda}] \right\} dS - \\ & \int_{S_{\sigma}} (p_{\lambda-} \ p_{\lambda}) \delta u_{\lambda} dS + \int_{S_u} (u_{\lambda-} \ u_{\lambda}) \delta p_{\lambda} dS, \end{aligned} \quad (25)$$

由(24)得

$$\delta\Gamma = 0, \quad (26)$$

当且仅当

$$1) \ W_{KJ}, \ \alpha_{\alpha} = \varepsilon_{\alpha}(u) \quad (\alpha = p, \dots, q) \quad (\text{在 } V_m \text{ 内}), \quad (27)$$

$$W_{KJ}, \ \varepsilon_{\beta} = -\sigma_{\beta} \quad (\beta = r, \dots, t) \quad (\text{在 } V_m \text{ 内}), \quad (28)$$

$$\mathfrak{E} = \mathfrak{E}(u) \quad (\beta = r, \dots, t) \quad (\text{在 } V_m \text{ 内}) \quad (29)$$

$$(F_{\lambda\lambda}, \sigma_{\lambda\mu}), \ + f_{\lambda} = 0 \quad (\text{在 } V_m \text{ 内}) \cdot \quad (30)$$

$$2) \ p \ d\lambda = - \ p \ u_{\lambda}, \ u \ u_{\lambda} = \ u \ u_{\lambda} \quad (\text{在 } S_m \text{ 上}) \cdot \quad (31)$$

$$3) \ u_{\lambda} = u_{\lambda} \quad (\text{在 } S_u \text{ 上}), \quad (32)$$

$$4) \ p_{\lambda} = p_{\lambda} \quad (\text{在 } S_{\sigma} \text{ 上}), \quad (33)$$

其中(27)是部分应力应变关系和应变位移关系的混合方程, (28), (29)是其余的应力应变关系和应变位移关系; (30)是关于  $V_m$  的平衡方程。

类似地, 由(25)可得

$$\delta\Pi = 0 \quad (34)$$

当且仅当(30)~(33)被满足, 并且

$$1) \ W_{KJ}, \ \mathfrak{E}_{\beta} = -\mathfrak{E}(u) \quad (\beta = r, \dots, t) \quad (\text{在 } V_m \text{ 内}), \quad (27)'$$

$$W_{KJ}, \ \varepsilon_{\alpha} = \alpha_{\alpha} \quad (\alpha = p, \dots, q) \quad (\text{在 } V_m \text{ 内}), \quad (28)'$$

$$\varepsilon_{\alpha} = \varepsilon_{\alpha}(u) \quad (\alpha = p, \dots, q) \quad (\text{在 } V_m \text{ 内}) \cdot \quad (29)'$$

注意到(22),  $n \ d\lambda = - \ n \ u_{\lambda}$ , 从(31a), 我们得

$$\begin{aligned} & [ \delta_{\lambda\lambda} + u \ u_{\lambda} ] \sigma \ \lambda_3 = \\ & [ \delta_{\lambda} \ \lambda + u \ u_{\lambda} ] \sigma \ \lambda_3 \quad (\text{在 } S_m \text{ 上}), \end{aligned} \quad (35)$$

由此, 对于小变形, (35)成为

$$\sigma \ \lambda_3 = \sigma \ \lambda_3 \quad (\text{在 } S_m \text{ 上}) \cdot \quad (36)$$

应力和位移的连续条件(31)在层合板的理论和应用中是很重要的<sup>[2~5]</sup>。

注意到  $[C_{ij}][H_{ij}] = I, Rr = I$ ; (7)分别等价于(8)和(9), 易见, 在(5), (7)~(11), (14)~(16)中只要有一个被满足, 则其余的全满足, 进而推出

$$\Pi + \Gamma = 0 \cdot \quad (37)$$

### 3 算 例

本文对铺层为  $[0^{\circ}/90^{\circ}/0^{\circ}]$  ( $b/a = 3$ ) 和  $[0^{\circ}/90^{\circ}/90^{\circ}/0^{\circ}]$  ( $b/a = 1, b$  和  $a$  分别为板面的长

和宽)、四边简支矩形层合板(各层厚度均相等)在上表面载荷  $p_3 = p_0 \sin(\pi x_1/a) \sin(\pi x_2/b)$  作用下,利用变分式  $\delta \Gamma(u, \sigma_i, \varepsilon_i, \varepsilon_5) = 0$ , 通过  $4 \times 4$  的网格的有限元法进行了数值计算,并求得了板的中点处的挠度,(利用了对称性只取板的四分之一进行了计算)• 并将  $a/h = 100$ , ( $h$  为板的厚度)的计算结果与已见到的结果比较如下:

i)  $[0^\circ/90^\circ/0^\circ]$  ( $b/a = 3$ )

本文:  $W = 0.513$ , 文[6]:  $W = 0.505$

其中  $W = \left\{ 100 E_2 h^3 / (p_0 a^4) \right\} W_0$ ,  $W_0$  为中点处的挠度•

ii)  $[0^\circ/90^\circ/90^\circ/0^\circ]$  ( $b/a = 1$ )

本文:  $W = 1.025$ , 文[6]:  $W = 1.003$ ,

其中  $W = \left\{ 100 E h^3 / (p_0 a^4) \right\} W_0$ ,  $W_0$  为中点处的挠度•

$$E = (\pi^4/12) \left\{ 4G_{12} + [E_1 + (1 + 2\nu)E_2] / (1 - \nu_{12}\nu_{21}) \right\}$$

在计算中,每层材料的弹性常数为

$$E_1/E_2 = 25, \quad G_{12}/E_2 = 0.5, \quad \nu_{12} = 0.25 \text{ (从而求得 } \nu_{21} = 0.01)$$

数值结果表明,本文提供的泛函是可靠的•

### [参 考 文 献]

- [1] Christensen R M. Mechanics of Composite Materials [M]. New York John Wilen & Sons, 1979, 153.
- [2] Spilker R L, Chou S C. Edge effects in symmetric composite laminates importance of satisfying the traction\_free\_edge condition[J]. J Compos Mater, 1980, 14: 2~ 20.
- [3] Bar Yoseph P, Pian T H H. Calculation of interlaminar stress concentration in composite laminates [J]. J Compos Mater, 1981, 15(3): 225~ 239.
- [4] Alturi E A Rotem Shmueli M. Free edge effect in angle ply laminates—a new three dimensional finite difference solution[J]. J Compos Mater, 1980, 14(1): 21~ 30.
- [5] Yamada M, Nemat\_Nasser S. Harmonic waves with arbitrary propagation direction in layered orthotropic elastic composites[J]. J Compos Mater, 1981, 15(1): 531~ 542.
- [6] Pagano N J, Exact solutions for rectangular bidirectional composites and sandwich plates[J]. J Compos Mater, 1970, 4(1): 20~ 34.

## Continual Conditions Between Layers for Laminated Plates

Zhang Shenxue

(Department of Mathematics, Jilin University, Changchun 130023, P R China)

**Abstract:** In this paper, various forms of functional on blending energy principles of composite laminated plates are given, which guarantee satisfied continual conditions of displacements and stress between layers, and then the reliability of the functional are proved by the computing example.

**Key words:** laminated plates; continual conditions between layers; blending energy; functional