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非线性弹性杆波动方程的显式精确解*

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摘要: 应用 sine-cosine 方法对非线性弹性杆波动方程进行了求解, 得到了该方程的一些新的周期波解和孤波解 (材料常数 n 为不等于 1 的常数). 对部分结果通过数学软件得到了了解的图像, 获得的结果有助于非线性弹性杆中孤波存在性问题的进一步研究.

关键词: sine-cosine 方法; 非线性弹性杆; 精确解; 孤波

中图分类号: O343; O175.2 **文献标志码:** A **DOI:** 10.21656/1000-0887.420245

Explicit Exact Solutions to the Wave Equation for Nonlinear Elastic Rods

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Abstract: The sine-cosine method was applied to the wave equation for nonlinear elastic rods, and some new periodic and solitary solutions to the equation were obtained (with material constant n other than 1). The graphs of some solutions were given through the math software. The results are helpful to further research on existence of solitary waves in the nonlinear elastic rods.

Key words: sine-cosine method; nonlinear elastic rod; exact solution; solitary wave

引 言

杆是重要的工程应用元件, 杆件中的波动特别是非线性波的研究日益引起科研人员的关注, 现已成为工程领域的一个研究热点^[1-6]. 在文献 [1] 中, 张善元和庄蔚建立了计入横向惯性效应任意形状截面的非线性弹性杆波动方程, 其形式为

* 收稿日期: 2021-08-17; 修订日期: 2021-11-25

基金项目: 国家自然科学基金 (11765017; 12065022; 12165018); 甘肃省重点人才项目 (2020RCXM100)

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引用格式: 郭鹏, 唐荣安, 孙小伟, 洪学仁, 石玉仁. 非线性弹性杆波动方程的显式精确解 [J]. 应用数学和力学, 2022, 43(8): 869-876.

$$\frac{\partial^2 u}{\partial t^2} - c_0^2 \left[1 + na_n \left(\frac{\partial u}{\partial x} \right)^{n-1} \right] \frac{\partial^2 u}{\partial x^2} - \frac{v^2 J_\rho}{s} \frac{\partial^4 u}{\partial t^2 \partial x^2} = 0, \quad (1)$$

其中, s 为杆截面积, J_ρ 为杆截面的极惯性矩, $c_0^2 = E/\rho$ 为线弹性纵波波速平方, E 为弹性模量, ρ 为杆的密度, ν 为 Poisson 比, a_n 和 n 均为材料常数. 取材料常数 $n = 2$, 通过变换将方程变换为 KdV 方程并用逆散射方法获得了 KdV 方程的孤波解, 证明了在材料常数 $n = 2$ 时, 非线性弹性杆中存在孤波. 很多学者对非线性弹性杆中的孤波问题进行了研究. 在文献 [7] 中, 胡伟鹏等取材料常数 $n = 2$, 用多辛方法对方程 (1) 进行了数值模拟, 讨论了非线性效应和几何弥散效应对孤波传播的影响. 在文献 [8] 中, Duan 和 Zhao 取材料常数 $n = 3$, 用约化摄动方法将方程 (1) 变换为非线性 Schrödinger 方程, 得到了方程的包络孤波解. 在文献 [9] 中, 郭鹏等取材料常数 $n \geq 2$, 用约化摄动方法将方程 (1) 变换为变形 KdV 方程. 在文献 [10] 中, 吕克璞等取材料常数 $n \geq 2$, 用约化摄动方法将方程 (1) 变换为非线性 Schrödinger 方程. 以上研究证明了在材料常数 $n \geq 2$ 时, 非线性弹性杆中存在孤波和包络孤波. 但是需要注意的是, 用约化摄动方法获得的解还是方程的近似解. 在文献 [11] 中, Kabir 取材料常数 $n = 2$, 用修正的 Kudryashov 方法、 (G'/G) 展开法和 \exp 函数法获得了方程 (1) 的精确解. 在文献 [12] 中, Çelik 等取材料常数 $n = 2$, 用 Lie 群分析方法和 F -展开法获得了方程 (1) 的精确解. 在文献 [13-14] 中, Li 等用平面动力系统方法研究了材料常数 $n = 2, 3, 5$ 以及大于 5 的奇数情况下的精确解.

非线性弹性杆波动方程的求解尽管已经有了上述研究工作, 但是对于方程 (1) 中材料常数 n 取不同数值时 (包括分数), 求解 (特别是精确解) 仍然是很困难的. 值得注意的是, 对于非线性方程的求解, 经过相关学者近几十年来的努力, 目前已经发展出了很多精确求解的方法^[15-36]. 在这些求解方法中, sine-cosine 方法是一种较为简便的方法. 这种方法的思路是将方程的解拟设为正弦或余弦函数形式, 然后将正弦或余弦函数形式的解代入原方程, 用待定系数法确定相应的常数, 从而获得方程的精确解. 本文应用 sine-cosine 方法^[29-33] 对非线性弹性杆波动方程进行了求解, 得到了该方程的一些新的周期波解和孤波解. 对材料常数 n 不等于 1 (包括材料常数 n 为不等于 1 的分数) 的非线性弹性杆中孤波的存在性给出了理论证明.

1 非线性弹性杆波动方程的求解

对方程 (1) 引入如下变换:

$$\xi = kx - ct, \quad (2)$$

其中 k, c 为待定常数. 将方程 (2) 代入方程 (1) 可得

$$\frac{\partial^2 u}{\partial \xi^2} - \frac{na_n c_0^2 k^3}{c^2 - c_0^2 k^2} \left(\frac{\partial u}{\partial \xi} \right)^{n-1} \frac{\partial^2 u}{\partial \xi^2} - \frac{v^2 J_\rho c^2 k^2}{s(c^2 - c_0^2 k^2)} \frac{\partial^4 u}{\partial \xi^4} = 0. \quad (3)$$

将方程 (3) 对 ξ 积分一次, 并取积分常数为零可得

$$\frac{\partial u}{\partial \xi} - \frac{a_n c_0^2 k^3}{c^2 - c_0^2 k^2} \left(\frac{\partial u}{\partial \xi} \right)^n - \frac{v^2 J_\rho c^2 k^2}{s(c^2 - c_0^2 k^2)} \frac{\partial^3 u}{\partial \xi^3} = 0. \quad (4)$$

令

$$w = \frac{\partial u}{\partial \xi}, \quad (5)$$

则方程 (4) 的形式可简化为

$$w - \alpha w^n - \beta \frac{\partial^2 w}{\partial \xi^2} = 0, \quad (6)$$

其中

$$\alpha = \frac{a_n c_0^2 k^3}{c^2 - c_0^2 k^2}, \beta = \frac{v^2 J_\rho c^2 k^2}{s(c^2 - c_0^2 k^2)}. \quad (7)$$

假设方程 (6) 的解为

$$w(\xi) = \lambda \sin^m(\mu \xi). \quad (8)$$

将方程 (8) 代入方程 (6) 可得

$$\lambda(1 + \beta m^2 \mu^2) \sin^m(\mu\xi) - \alpha \lambda^n \sin^{nm}(\mu\xi) - \beta(m-1)m\lambda\mu^2 \sin^{m-2}(\mu\xi) = 0. \quad (9)$$

为使方程 (9) 有非零解, 需要平衡其中第二和第三项正弦函数的指数, 使得

$$nm = m - 2. \quad (10)$$

从方程 (10) 可以很方便地确定出 $m = -2/(n-1)$. 这样材料常数 n 只能取为不等于 1 的常数. 将 m 代入方程 (9) 并合并正弦函数的同次项, 可得

$$\lambda \left[1 + \frac{4\beta\mu^2}{(1-n)^2} \right] \sin^{-\frac{2}{n-1}}(\mu\xi) - \left[\alpha\lambda^n + \frac{2\beta(1+n)\lambda\mu^2}{(1-n)^2} \right] \sin^{-\frac{2n}{n-1}}(\mu\xi) = 0. \quad (11)$$

令方程 (11) 中正弦函数前的系数为零, 可得

$$\begin{cases} \lambda \left[1 + \frac{4\beta\mu^2}{(1-n)^2} \right] = 0, \\ \alpha\lambda^n + \frac{2\beta(1+n)\lambda\mu^2}{(1-n)^2} = 0. \end{cases} \quad (12)$$

由方程 (12) 解得

$$\lambda = \frac{n-1}{2\alpha} \sqrt{\frac{n+1}{2\alpha}}, \quad \mu = \pm(n-1) \sqrt{-\frac{1}{4\beta}}. \quad (13)$$

将方程 (13) 及 α, β 代入方程 (8), 可得

$$w_1 = \frac{n-1}{\sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}}} \csc^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (n-1)\xi \right], \quad (14)$$

$$w_2 = \frac{n-1}{\sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}}} \csc^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (n-1)\xi \right]. \quad (15)$$

同理, 若假设方程 (6) 的解为

$$w(\xi) = \lambda \cos^m(\mu\xi). \quad (16)$$

将方程 (16) 代入方程 (6), 得到的解为

$$w_3 = \frac{n-1}{\sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}}} \sec^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (n-1)\xi \right], \quad (17)$$

$$w_4 = \frac{n-1}{\sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}}} \sec^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (n-1)\xi \right]. \quad (18)$$

当 $c_0^2 k^2 - c^2 < 0$ 时, 用双曲函数与三角函数之间的关系式

$$\csc(ix) = -i \operatorname{csch} x, \quad \sec(ix) = \operatorname{sech} x, \quad (19)$$

可将方程 (14)、(15)、(17)、(18) 化为

$$w_5 = \frac{n-1}{\sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}}} (-i)^{\frac{2}{n-1}} \operatorname{csch}^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (n-1)\xi \right], \quad (20)$$

$$w_6 = \frac{n-1}{\sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}}} (-i)^{\frac{2}{n-1}} \operatorname{csch}^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (n-1)\xi \right], \quad (21)$$

$$w_7 = \frac{n-1}{\sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}}} \operatorname{sech}^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (n-1)\xi \right], \quad (22)$$

$$w_8 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} \operatorname{sech}^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi. \quad (23)$$

根据方程(5),将方程(14)、(15)、(17)、(18)对 ξ 积分一次,可得

$$u_1 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} \operatorname{csc}^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi, \quad (24)$$

$$u_2 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} \operatorname{csc}^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi, \quad (25)$$

$$u_3 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} \operatorname{sec}^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi, \quad (26)$$

$$u_4 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} \operatorname{sec}^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi. \quad (27)$$

方程(24)~(27)即为非线性弹性杆波动方程的三角函数周期波解.

同理,将方程(20)~(23)对 ξ 积分一次,可得

$$u_5 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} (-i)^{\frac{2}{n-1}} \operatorname{csch}^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi, \quad (28)$$

$$u_6 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} (-i)^{\frac{2}{n-1}} \operatorname{csch}^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi, \quad (29)$$

$$u_7 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} (-i)^{\frac{2}{n-1}} \operatorname{sech}^{\frac{2}{n-1}} \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi, \quad (30)$$

$$u_8 = \int^{n-1} \sqrt{\frac{(n+1)(c^2 - c_0^2 k^2)}{2a_n c_0^2 k^3}} (-i)^{\frac{2}{n-1}} \operatorname{sech}^{\frac{2}{n-1}} \left[-\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4v^2 J_\rho c^2 k^2}} (n-1)\xi \right] d\xi. \quad (31)$$

方程(28)~(31)即为非线性弹性杆波动方程的孤波解.

有了以上精确解的表达式,只要确定了材料常数 n ,即可通过积分获得非线性弹性杆波动方程的精确解.

例如,当材料常数 $n = 1.5$ 时,由方程(24)~(31)可得

$$u_1 = -\frac{(2.5)^2 (c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4v^2 J_\rho c^2 k^2}{s(c_0^2 k^2 - c^2)}} \cot \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[2 + \operatorname{csc}^2 \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (32)$$

$$u_2 = -\frac{(2.5)^2 (c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4v^2 J_\rho c^2 k^2}{s(c_0^2 k^2 - c^2)}} \cot \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[2 + \operatorname{csc}^2 \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (33)$$

$$u_3 = -\frac{(2.5)^2 (c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4v^2 J_\rho c^2 k^2}{s(c_0^2 k^2 - c^2)}} \tan \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[2 + \operatorname{sec}^2 \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (34)$$

$$u_4 = -\frac{(2.5)^2 (c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4v^2 J_\rho c^2 k^2}{s(c_0^2 k^2 - c^2)}} \tan \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[2 + \operatorname{sec}^2 \left[0.5 \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4v^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (35)$$

$$u_5 = -\frac{(2.5)^2(c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4\nu^2 J_\rho c^2 k^2}{s(c^2 - c_0^2 k^2)}} \coth \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[-2 + \operatorname{csch}^2 \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (36)$$

$$u_6 = -\frac{(2.5)^2(c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4\nu^2 J_\rho c^2 k^2}{s(c^2 - c_0^2 k^2)}} \coth \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[-2 + \operatorname{csch}^2 \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (37)$$

$$u_7 = -\frac{(2.5)^2(c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4\nu^2 J_\rho c^2 k^2}{s(c^2 - c_0^2 k^2)}} \tanh \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[2 + \operatorname{sech}^2 \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (38)$$

$$u_8 = -\frac{(2.5)^2(c^2 - c_0^2 k^2)^2}{6a_n^2 c_0^4 k^6} \sqrt{\frac{4\nu^2 J_\rho c^2 k^2}{s(c^2 - c_0^2 k^2)}} \tanh \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \times \left[2 + \operatorname{sech}^2 \left[0.5 \sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] \right] + C, \quad (39)$$

其中, C 为积分常数.

当材料常数 $n = 3$ 时, 由方程 (24)~(31) 可得

$$u_1 = \sqrt{-\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} \ln \tan \left[\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] + C, \quad (40)$$

$$u_2 = \sqrt{-\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} \ln \cot \left[\sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] + C, \quad (41)$$

$$u_3 = \sqrt{-\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} \ln \frac{1 + \tan \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct)}{1 - \tan \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct)} + C, \quad (42)$$

$$u_4 = \sqrt{-\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} \ln \frac{1 + \tan \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct)}{1 - \tan \sqrt{\frac{s(c_0^2 k^2 - c^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct)} + C, \quad (43)$$

$$u_5 = \sqrt{\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} (-i) \ln \tanh \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] + C, \quad (44)$$

$$u_6 = \sqrt{\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} (-i) \ln \coth \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] + C, \quad (45)$$

$$u_7 = 2 \sqrt{\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} (-i) \arctan \tanh \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] + C, \quad (46)$$

$$u_8 = 2 \sqrt{\frac{8\nu^2 J_\rho c^2}{a_n c_0^2 k s}} (-i) \arctan \tanh \left[\sqrt{\frac{s(c^2 - c_0^2 k^2)}{4\nu^2 J_\rho c^2 k^2}} (kx - ct) \right] + C, \quad (47)$$

其中, C 为积分常数. 当材料常数 $n = 3$ 时, 在文献 [13] 中得到了方程 (1) 如下形式的精确解:

$$u = \pm \left(\frac{4\sqrt{\delta}}{\gamma} \right) \arctan(e^{\sqrt{\delta}(x-\omega t)}), \quad (48)$$

$$u = \pm \left(\sqrt{\frac{2}{-\gamma}} \right) \ln \cosh \left[\sqrt{\frac{-\delta}{2}} (x - \omega t) \right], \quad (49)$$

其中, $\delta = \frac{s(c^2 - c_0^2)}{c^2 v^2 J_p}$, $\gamma = \frac{s c_0^2 a_n}{c^2 v^2 J_p}$, ω 为波速. 通过比较可以看出, 本文所得结果为形式不同的新解.

由材料常数 n 取其他不同的数值, 我们还可以获得更多新的精确解, 此处不再一一列出.

2 结果与讨论

应用 sine-cosine 方法对非线性弹性杆波动方程进行了求解, 得到了一些新的周期波解和孤波解. 根据获得的精确解, 还可以很方便地绘制精确解的图像. 例如图 1 绘制了当 $c = 1.5$, $k = 1$, $a_n = 0.7$, $c_0 = 100$, $v = 0.4$, $J_p = 0.3$, $s = 3.5$, $C = 0$ 时, 方程 (34) 的三角函数周期波解; 图 2 绘制了当 $c = 100$, $k = 0.5$, $a_n = 0.7$, $c_0 = 30$, $v = 0.2$, $J_p = 0.1$, $s = 10$, $C = 0$ 时, 方程 (38) 的孤波解. 所获得的非线性弹性杆波动方程的新精确解 (特别是孤波解) 对材料常数 n 不等于 1 的非线性弹性杆中孤波的存在性给出了理论证明 (包括材料常数 n 为不等于 1 的分数), 有助于该问题的进一步研究. 从求解非线性弹性杆波动方程的过程可以看出, sine-cosine 方法是一种简便、有效的方法. 在文献 [32-33] 中, 采用 sine-cosine 方法同样简便地获得了 KdV 方程、修正 KdV 方程、广义 KdV 方程、Boussinesq 方程、RLW 方程、BBM 方程、Phi-4 方程、修正 Degasperis-Procesi 方程、修正 Camassa-Holm 方程的精确解. 这种方法还可用于其他非线性方程或方程组的求解.

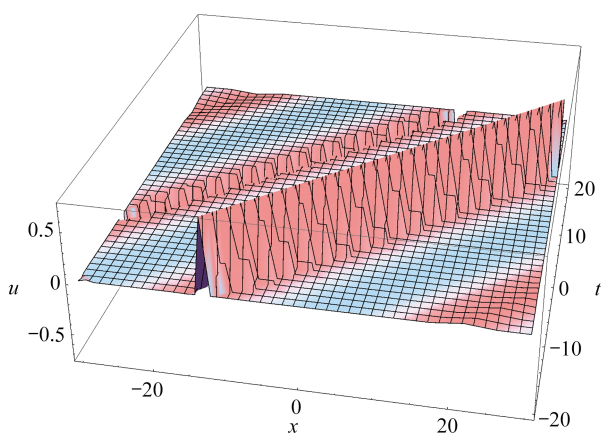


图 1 方程 (34) 解的图像

Fig. 1 Graphical representation of solution (34)

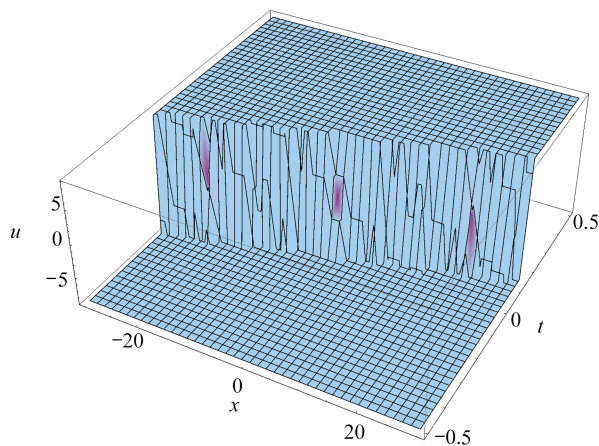


图 2 方程 (38) 解的图像

Fig. 2 Graphical representation of solution (38)

3 结束语

非线性弹性杆波动方程的求解尽管已经有了相关参考文献及本文所开展的上述研究工作, 但是对于方程 (1) 中材料常数取不同数值时 (包括分数) 的求解仍然是很困难的. 继续研究新的求解方法与获得同类型方程更丰富的解, 值得研究工作者们进行更深入的钻研与探索.

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